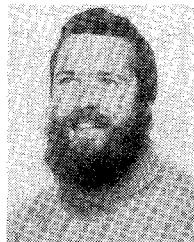


- [11] H. F. Taylor, "Power loss at directional change in dielectric waveguides," *Appl. Opt.*, vol. 13, pp. 642-647, Mar. 1974.
- [12] J. E. Goell and R. D. Standley, "Sputtered glass waveguide for integrated optical circuits," *Bell Syst. Tech J.*, vol. 48, pp. 3445-3448, Dec. 1969.
- [13] J. E. Goell, "Electron-resist fabrication of bends and couplers for integrated optical circuits," *Appl. Opt.*, vol. 12, pp. 729-736, Apr. 1973.
- [14] J. E. Goell, "Rib waveguide for integrated optical circuits," *Appl. Opt.*, vol. 12, pp. 2797-2798, Dec. 1973.
- [15] H. Furuta, H. Noda, and A. Ihaya, "Novel optical waveguide for integrated optics," *Appl. Opt.*, vol. 13, pp. 322-326, Feb. 1974.
- [16] L. D. Hutcheson, I. A. White, and J. J. Burke, "Losses in diffused LiNbO_3 waveguides caused by directional changes," in *Dig. of Top. Meet. Integ. and Guided-Wave Opt.*, Incline Village, NV, Jan. 1980, paper WB2.
- [17] R. G. Walker, "The design of ring-resonators for integrated optics using silver ion-exchanged waveguides," Ph.D. thesis, University of Glasgow, 1981.

- [18] J. T. Boyd, "Theory of parametric oscillation phase matched in GaAs thin-film waveguides," *IEEE J. Quantum Electron.*, vol. QE-8, pp. 788-796, Oct. 1972.



Michael W. Austin was born in Melbourne, Australia in 1955. He received the B.E. degree with distinction in communication engineering in 1977 and the Masters degree in electronic engineering in 1980, both from the Royal Melbourne Institute of Technology, Melbourne, Australia.

Since late 1979 he has been employed as a Research Fellow at the British Telecom Research Laboratories, Ipswich, England, where he is currently working towards the Ph.D. in the field of integrated optics in III-V semiconductors.

Mr. Austin is a member of IEEE Australia.

CO₂ Laser Annealing of Si₃N₄, Nb₂O₅, and Ta₂O₅ Thin-Film Optical Waveguides to Achieve Scattering Loss Reduction

SUBHADRA DUTTA, STUDENT MEMBER, IEEE, HOWARD E. JACKSON, J. T. BOYD, SENIOR MEMBER, IEEE, RICHARD L. DAVIS, AND FRED S. HICKERNELL, SENIOR MEMBER, IEEE

Abstract—Significant reductions in the optical scattering losses of Si₃N₄, Nb₂O₅, and Ta₂O₅ waveguides fabricated on SiO₂/Si substrates have been measured following CO₂ laser annealing. The largest improvements were observed for Si₃N₄ waveguides, where waveguide attenuation values of about 6.0 dB/cm before laser annealing were reduced to as low as 0.1 dB/cm afterwards. An improvement of more than an order of magnitude was obtained for a Nb₂O₅ waveguide upon laser annealing, the attenuation coefficient decreasing from 7.4 to 0.6 dB/cm. In the case of one Nb₂O₅ waveguide no improvement was obtained upon laser annealing. The attenuation coefficient of a reactively sputtered Ta₂O₅ waveguide was found to decrease from 1.3 dB/cm before laser annealing to 0.4 dB/cm afterwards. In the case of a thermally oxidized Ta₂O₅ waveguide a small initial improvement in waveguide attenuation was followed by degradation upon further laser annealing.

Manuscript received June 24, 1981; revised September 17, 1981. This work was supported in part by the Air Force Office of Scientific Research and the National Science Foundation.

S. Dutta was with the Department of Physics, University of Cincinnati, Cincinnati, OH 45221. She is now with the Westinghouse Research and Development Center, Pittsburgh, PA 15235.

H. E. Jackson is with the Department of Physics, University of Cincinnati, Cincinnati, OH 45221.

J. T. Boyd is with the Solid State Electronics Laboratory, Department of Electrical and Computer Engineering, University of Cincinnati, Cincinnati, OH 45221.

R. L. Davis and F. S. Hickernell are with the Motorola Government Electronics Division, Scottsdale, AZ 85252.

I. INTRODUCTION

SCATTERING of light propagating in thin-film optical waveguides represents a loss mechanism and, concommittantly, a limitation on the dynamic range obtainable in integrated optical signal processing devices [1]. An order of magnitude reduction in waveguide scattering loss, for instance, leads to an approximately 10 dB increase in dynamic range [1]. In addition, a lower loss waveguide requires a lower power light source, an advantage in some applications. Reduction of scattering in thin-film optical waveguides is thus desirable, as it would make integrated optical signal processing devices more competitive with other approaches.

We have previously reported CO₂ laser annealing that dramatically reduces scattering losses in Corning 7059 glass [2], [3] and ZnO [4] thin-film waveguides fabricated on thermally oxidized silicon substrates. Losses as low as 0.05 dB/cm for Corning 7059 glass waveguides and 0.01 dB/cm for ZnO waveguides have been achieved by this technique, a factor of 50-100 times lower than the best results previously reported for these two waveguide materials [5]-[10]. In this paper we wish to report the success of CO₂ laser annealing in reducing scattering losses in Si₃N₄, Nb₂O₅, and Ta₂O₅ waveguides, all fabricated on thermally oxidized silicon substrates.

The lowest loss reported for any thin-film optical waveguide

prior to our recently reported laser annealing results [3], [4] had been 0.1 dB/cm for a Si₃N₄ waveguide, [11]. This makes Si₃N₄ a promising waveguide material despite difficulties experienced in fabricating good quality films [11], [12]. We have reduced the waveguide attenuation of Si₃N₄ waveguides from values of about 6.0 dB/cm before laser annealing to as low as 0.1 dB/cm after laser annealing. Our success in reducing the waveguide loss by laser annealing implies that the film deposition conditions are no longer as critical to the fabrication of low-loss Si₃N₄ waveguides. Nb₂O₅ has been used as a low-loss waveguide material for several years [5], [13], and its high refractive index makes it potentially useful for waveguide lens applications as well [14]. We have reduced the scattering loss of a Nb₂O₅ waveguide from 7.4 to 0.6 dB/cm by laser annealing. Ta₂O₅ has long been considered an attractive material for integrated optical device applications, and it is widely used for thin-film waveguide lenses because of its high refractive index [15]–[17]. We have reduced the waveguide attenuation of an RF sputtered Ta₂O₅ waveguide from 1.2 dB/cm before laser annealing to 0.4 dB/cm after laser annealing.

In what follows we first describe, in Section II, the fabrication of the Si₃N₄ films by low pressure chemical vapor deposition, the Nb₂O₅ films by RF sputtering, and the Ta₂O₅ films by both reactive sputtering and thermal oxidation of sputtered tantalum films. We then detail, in Section III, the laser annealing conditions for each type of waveguide and our technique for the measurement of low values of waveguide loss. Here we include a discussion of the errors introduced by the statistical fluctuations inherent in this type of measurement. In Section IV we present and discuss the optical waveguide attenuation data obtained before and after laser annealing, and the refractive index changes induced by laser annealing. Finally, in Section V, we summarize our conclusions for the three different types of waveguides.

II. WAVEGUIDE FABRICATION

The Si₃N₄ waveguides were deposited on 3 in diameter silicon substrates, each of which had a ~1.0 μ m thick surface layer of SiO₂ formed by thermal oxidation at 1100°C for 2.5 h in a steam environment. The Si₃N₄ was deposited using a low pressure chemical vapor deposition (LPCVD) furnace at a temperature of 800°C. The LPCVD process used was based on the thermal reaction of dichlorosilane (H₂SiCl₂) and anhydrous ammonia (NH₃) in a hot wall diffusion furnace at a pressure of approximately 350 mtorr. The deposition rate was 0.32 \AA/s . Film thicknesses of 0.280 μ m for sample Si₃N₄-1 and 0.175 μ m for sample Si₃N₄-2 were obtained. The films were uniform in thickness to within ± 1 percent as determined by optical measurements. There were no microfractures evident in the films, as seen by prism-coupling He-Ne laser light into the waveguides.

The Nb₂O₅ waveguides, samples Nb₂O₅-1 and Nb₂O₅-2, were formed by reactively sputtering niobium metal in an argon-oxygen atmosphere. The optical properties of these films are strongly dependent upon the conditions under which the films are grown [18]. This is illustrated by the two samples used in this study. Two-inch diameter silicon wafers with ~1.0 μ m thick layers of SiO₂ formed on their surfaces were

used as substrates for the optical waveguides. Both Nb₂O₅ films were deposited in an RF diode sputtering machine with bias sputtering capability. Both samples were bias sputtered with a total chamber gas pressure of 4.0 mtorr and a ratio of argon to oxygen of 3:1. The total RF power of sample Nb₂O₅-1 was 150 W, with the power divided approximately 4:1 between sputtering target and substrate, resulting in a deposition rate of 6.0 \AA/min . This led to bias voltages of -810 V on the target and -190 V on the substrate. The film thickness was measured to be about 0.742 μ m. For sample Nb₂O₅-2, the total power was 50 W, with a 6:1 power ratio. The bias voltages were -450 and -75 V for target and substrate, respectively. This resulted in a deposition rate of 4.3 \AA/min and a film thickness of 0.518 μ m.

Of the two Ta₂O₅ waveguides, the first, sample Ta₂O₅-1, was formed by the oxidation of sputtered tantalum metal, and the second, sample Ta₂O₅-2, was formed directly by reactive sputtering. In both cases the films were deposited on thermally oxidized silicon substrates in an RF magnetron system using a high purity tantalum target. The tantalum metal was deposited at a background pressure of 10.0 mtorr with an argon:oxygen ratio of 97:3 using 500 W applied power. The accumulation rate was 1.5 \AA/s . The directly sputtered Ta₂O₅ was also deposited with a background gas pressure of 10.0 mtorr but with oxygen increased to a 90:10 ratio and applied power of 700 W. The accumulation rate was 0.3 \AA/s . The relative amount of oxygen in the ionized gas accounts for the formation of the metal or the oxide, with the transition occurring at an argon:oxygen ratio of approximately 95:5. The addition of slight amounts of oxygen in the metal deposition result in a better and less process sensitive Ta₂O₅ optical waveguide after oxidation [19]. The resistivities of the tantalum metal film and the reactively sputtered Ta₂O₅ film were $4.5 \times 10^{-4} \Omega \cdot \text{cm}$ and $<10^{13} \Omega \cdot \text{cm}$, respectively. The tantalum metal film was oxidized at 550°C for 24 h in an open tube furnace with an oxygen flow rate of approximately 2500 cm^3/min . This brought the resistivity to above $10^{12} \Omega \cdot \text{cm}$. The dielectric properties of the thermally oxidized and reactively sputtered films were measured as capacitors at 1.0 kHz. They had dielectric constants of 25.0 and 19.0, and dissipation factors of 0.0008 and 0.0002, respectively. The variations in the dielectric properties are attributed to stoichiometry and density differences of the films. The thermally oxidized film was 0.510 μ m thick, and the reactively sputtered film was 0.445 μ m thick.

III. LASER ANNEALING AND WAVEGUIDE CHARACTERIZATION TECHNIQUES

Laser annealing was carried out using a 50 W, CW CO₂ laser, a 6.4 cm focusing lens, and a variable frequency beam scanner. The beam was scanned horizontally across a distance of 1 cm at a rate of 1.0 cm/s, and the sample, which was mounted vertically on a motorized translation stage, was translated at the end of each horizontal beam scan in vertical increments ranging from 150–200 μ m. A laser power density of $1.6 \times 10^5 \text{ W/cm}^2$ associated with a spot size of ~150 μ m was used to anneal the Si₃N₄ waveguides. Laser annealing was carried out at power densities varying from $3.0 \times 10^3 \text{ W/cm}^2$ to

1.0×10^5 W/cm², associated with spot sizes varying from 900 to ~ 150 μm , for the Nb_2O_5 waveguides. The Ta_2O_5 waveguides were laser annealed at power densities varying from 9.3×10^2 W/cm² to 2.0×10^5 W/cm², associated with spot sizes varying from 1.5 mm-150 μm .

Careful measurements of the out-of-plane scattering loss of each of these waveguides were made both before and after laser annealing. He-Ne laser light was prism-coupled into each waveguide and the intensity of the light scattered out of the waveguide was measured as a function of distance along the axis of propagation to determine the attenuation coefficient. The measurement was carried out with a scanning photometric microscope incorporating a 50 μm aperture optical fiber probe within the microscope eyepiece and a 10 X microscope objective. This allowed precise positioning of the probe with respect to the waveguide surface. For each measurement along the axis of propagation, the microscope was defocused by 3 mm, the fiber probe was scanned a distance of 1 mm transverse to the axis at a rate of 1 mm/min, and an integrating digital voltmeter was used to integrate the signal over the scan [3]. Defocusing the microscope resulted in a spatial smoothing of the light scattered from isolated waveguide defects, thus smoothing the fluctuations in the data caused by the uneven distribution of scattering centers. A schematic of the experimental arrangement is given in Fig. 1. Defocusing, scanning, and integrating in this manner provided an appropriate average of the scattered light intensity. In addition, for the low-loss waveguides the laser power was monitored with a precision of ± 0.05 percent during the measurement, and the readings were normalized accordingly. Using these averaging techniques has allowed consistent measurement of waveguide attenuation values as low as 0.01 dB/cm [3].

Although we have reduced considerably the errors due to statistical fluctuations in the density of scattering centers in each waveguide by our careful averaging procedure, we have not eliminated them completely. In the case of waveguides which have losses greater than 1.0 dB/cm this statistical error representing the deviation of the data from the least squares fit never exceeds 10 percent. However, for lower-loss waveguides this uncertainty becomes considerably larger because the same deviation of a data point from the best straight line on a semilog plot of the scattered intensity versus distance will have a greater effect on a shallow slope than it will on a steep slope. For example, in Fig. 3, curve (a), omitting the fifth data point, which deviates significantly from the best straight line, changes the slope by no more than 0.6 percent. In Fig. 3, curve (c), however, omitting the second data point, whose deviation from the best straight line is a third of the deviation considered in curve (a), changes the slope from 0.6 to 0.1 dB/cm. The statistical uncertainty for each best fit is tabulated with the optical attenuation data in Tables I, II, and III. We have calculated the uncertainty in each best fit by taking an arithmetic average of the deviations of the data points from the best fit and using this to calculate a maximum and a minimum slope. The maximum slope is calculated from the straight line formed by adding the average deviation to the left-hand extreme of the best fit straight line and subtracting it from

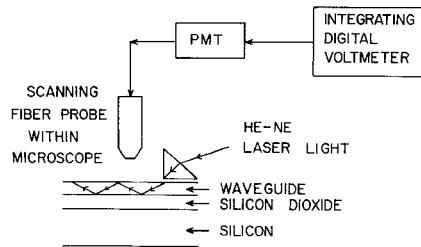


Fig. 1. Schematic of the experimental arrangement for measurement of out-of-plane waveguide scattering loss.

the right-hand extreme. The minimum slope is calculated by the reverse procedure. The uncertainty, then, is given by

$$\frac{\text{maximum slope} - \text{minimum slope}}{2\sqrt{n}}$$

where n is the number of data points. This method yields a slightly smaller uncertainty than that calculated using a similar formula for a least-squares fit error [20]. We wish to emphasize that this uncertainty is a result of the statistical distribution of scattering centers and not due to a lack of measurement precision. If we were to measure the waveguide attenuation along several different propagation axes across the entire width of the waveguide, we would expect that the differences in attenuation would be less than the statistical error in each best fit. This indicates that the fluctuation in the density of scattering centers is not a random error, so the error in the attenuation should be smaller than the statistical uncertainty tabulated, which is calculated assuming a random Gaussian distribution.

IV. RESULTS AND DISCUSSION

A. Si_3N_4 Waveguides

The optical attenuation data for the two Si_3N_4 waveguides, samples Si_3N_4 -1 and Si_3N_4 -2, are listed in Table I. Both samples were laser annealed at a power density of 1.6×10^5 W/cm². In sample Si_3N_4 -1, the waveguide attenuation decreased dramatically from 6.0 dB/cm before laser annealing to 0.1 dB/cm afterwards, a value equal to the lowest value previously reported for Si_3N_4 waveguides [11]. Scattering loss measurements for sample Si_3N_4 -1 before and after laser annealing are presented in Fig. 2(a). The waveguide attenuation in sample Si_3N_4 -2 also showed a significant decrease, from 3.1 dB/cm before laser annealing to 0.2 dB/cm afterwards, as shown in Fig. 2(b).

The refractive index of these Si_3N_4 films as determined by the measurement of TE and TM waveguide modes before laser annealing was 2.017 at the He-Ne laser wavelength ($\lambda = 6328$ Å). This value of index compares favorably with sputtered Si_3N_4 films formed by reactive sputtering [12]. The optical loss of LPCVD films was lower than that which had been achieved with the sputtered films. After laser annealing, the angles at which prism coupling occurred remained the same as before, indicating that the refractive index and thickness of neither of the films had changed upon laser annealing [21].

At the CO_2 laser wavelength, $\lambda = 10.6$ μm , Si_3N_4 has a very

TABLE I
EFFECT OF LASER ANNEALING ON Si₃N₄ WAVEGUIDE ATTENUATION

Sample Number	Si ₃ N ₄ Film Thickness (μm)	Loss Before Laser Annealing (dB/cm)	Loss After Laser Annealing (dB/cm)	Annealing Power Density (W/cm ²)
Si-1	0.280	6.0 ± 0.52	0.1 ± 0.5	1.6 × 10 ⁵
Si-2	0.175	3.1 ± 0.32	0.2 ± 0.8	1.6 × 10 ⁵

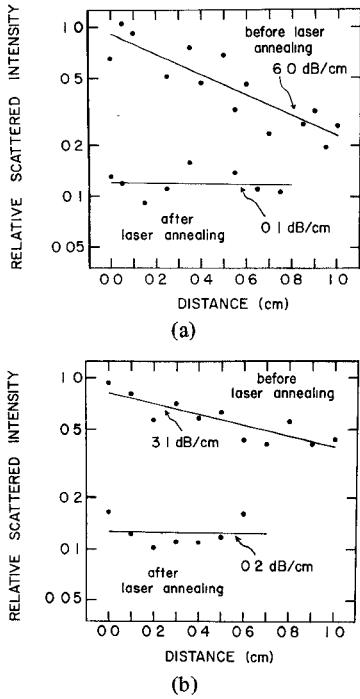


Fig. 2. Si₃N₄. Logarithm of the relative scattered intensity is plotted versus distance along the waveguide. (a) The upper and lower straight lines represent best fits to the data obtained from sample Si₃N₄-1 before and after laser annealing, respectively. (b) The upper and lower straight lines represent best fits to the data obtained from sample Si₃N₄-2 before and after laser annealing, respectively.

large absorption coefficient of $4.51 \times 10^4 \text{ cm}^{-1}$ [22]. This implies that although the Si₃N₄ film thicknesses were small (0.280 and 0.175 μm) the radiation absorbed generated sufficient heat to anneal out defects in these films. Thus, laser annealing resulted in significantly improved optical waveguides.

B. Nb₂O₅ Waveguides

Optical waveguide attenuation data for the two Nb₂O₅ waveguides, samples Nb₂O₅-1 and Nb₂O₅-2, are collected in Table II. For sample Nb₂O₅-1, a waveguide attenuation of 7.4 dB/cm was measured before laser annealing, as shown in Fig. 3, curve (a). This sample was initially laser annealed at a power density of $3.0 \times 10^3 \text{ W/cm}^2$, resulting in a lower attenuation of 3.6 dB/cm, as shown in Fig. 3, curve (b). A second laser anneal was carried out on sample Nb₂O₅-1 at a power density of $1.0 \times 10^5 \text{ W/cm}^2$. Using this higher power density for the second laser anneal, we observed a dramatic

reduction of the attenuation coefficient to 0.6 dB/cm, as shown in Fig. 2, curve (c). This value of loss is comparable to 0.5 dB/cm, the best result previously obtained for RF sputtered Nb₂O₅ waveguides [18].

The experiment was repeated on sample Nb₂O₅-2, which was a very high-loss waveguide. The laser light prism-coupled into this waveguide was attenuated so strongly that the waveguiding light streak was too short to allow any measurements of the attenuation coefficient to be made. This sample was laser annealed at $0.65 \times 10^5 \text{ W/cm}^2$. This did not result in any noticeable improvement of the waveguide, however. Visible damage occurred when a slightly higher laser power density of $1.0 \times 10^5 \text{ W/cm}^2$ was used.

The refractive indexes of the Nb₂O₅ films were measured before and after laser annealing and were found to remain essentially unchanged. At $\lambda = 6328 \text{ \AA}$ the refractive indexes were measured to be 2.186 for sample Nb₂O₅-1 and 2.286 for sample Nb₂O₅-2.

Nb₂O₅ has a fairly broad absorption band peaking at $\lambda = 11.8 \mu\text{m}$ [23], so that there is significant absorption at $\lambda = 10.6 \mu\text{m}$. Sample Nb₂O₅-1, which was a reasonably good waveguide to start with, showed moderate improvement after laser annealing at a power density of $3.0 \times 10^3 \text{ W/cm}^2$. When it was reannealed at the much higher power density of $1.0 \times 10^5 \text{ W/cm}^2$, the improvement was much more dramatic. A slightly higher power density of $1.3 \times 10^5 \text{ W/cm}^2$ caused visible damage, indicating that $1.0 \times 10^5 \text{ W/cm}^2$ was the optimum annealing power density for this waveguide. Sample Nb₂O₅-2, however, could not be improved by laser annealing.

C. Ta₂O₅ Waveguides

Table III displays optical attenuation data for the two Ta₂O₅ waveguides, samples Ta₂O₅-1 and Ta₂O₅-2, the thermally oxidized and reactively sputtered waveguides, respectively. A waveguide attenuation of 7.4 dB/cm was measured for sample Ta₂O₅-1 before laser annealing. Initially, laser annealing was tried at a power density of $3.0 \times 10^3 \text{ W/cm}^2$, resulting in a lower attenuation coefficient of 5.8 dB/cm. A second laser anneal was carried out at $1.3 \times 10^5 \text{ W/cm}^2$. The waveguide attenuation was left unchanged by this second anneal, a value of 5.7 dB/cm being measured afterwards. A higher power density of $2.0 \times 10^5 \text{ W/cm}^2$ was used for a third laser anneal. This caused a degradation of the waveguide attenuation to 9.5 dB/cm.

The reactively sputtered waveguide, sample Ta₂O₅-2, was of better optical quality than sample Ta₂O₅-1, and an attenuation of 1.3 dB/cm was measured before laser annealing, as shown

TABLE II
EFFECT OF LASER ANNEALING ON Nb_2O_5 WAVEGUIDE ATTENUATION

Sample Number	Nb_2O_5 Film Thickness (μm)	Loss Before Laser Annealing (dB/cm)	Power Density	Loss After First Anneal (dB/cm)	Power Density	Loss After Second Anneal (dB/cm)
			of First Anneal (W/cm^2)			
Nb-1	0.742	7.4 ± 0.63	3.0×10^3	3.6 ± 0.55	1.0×10^5	0.6 ± 0.26
Nb-2	0.518	100	0.65×10^5	100		

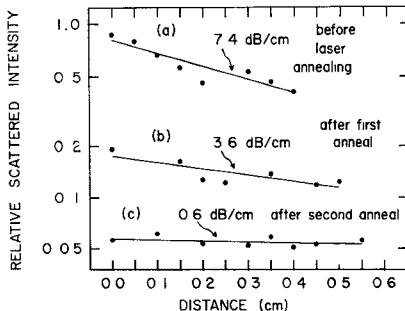


Fig. 3. Nb_2O_5 . Logarithm of the relative scattered intensity is plotted versus distance along the waveguide. The straight lines (a), (b), and (c) represent best fits to the data obtained from sample Nb_2O_5 -1 before laser annealing, after laser annealing at a power density of $3.0 \times 10^3 \text{ W/cm}^2$, and after a second laser anneal at a higher power density of $1.0 \times 10^5 \text{ W/cm}^2$, respectively.

TABLE III
EFFECT OF LASER ANNEALING ON Ta_2O_5 WAVEGUIDE ATTENUATION

Sample Number	Ta_2O_5 Film Thickness (μm)	Loss Before Laser Anneal- ing (dB/cm)	Power Density	Loss After First Anneal (dB/cm)	Power Density	Loss After Second Anneal (dB/cm)	Power Density	Loss After Third Anneal (dB/cm)
			of First Anneal (W/cm^2)				of Second Anneal (W/cm^2)	
Ta-1	0.45	7.4 ± 0.63	3.0×10^3	5.8 ± 0.50	1.3×10^5	5.7 ± 0.77	2.0×10^5	9.5 ± 0.58
Ta-2	0.45	1.3 ± 0.18	9.3×10^2	0.4 ± 0.28				

in Fig. 4, curve (a). Sample Ta_2O_5 -2 was laser annealed at a power density of $9.3 \times 10^2 \text{ W/cm}^2$. The waveguide attenuation decreased to 0.4 dB/cm after laser annealing, as shown in Fig. 4, curve (b). This compares favorably with the best result previously measured for Ta_2O_5 waveguides [15].

The refractive indexes of the two Ta_2O_5 films were measured using TE and TM waveguide modes both before and after laser annealing. At $\lambda = 6328 \text{ \AA}$, sample Ta_2O_5 -1 had an average index of 2.214 before laser annealing, which decreased to 2.057 after laser annealing. The refractive index of sample Ta_2O_5 -2 decreased from 2.120 before laser annealing to 2.106 after laser annealing. A decrease in index has also been observed after conventional thermal annealing in the 500 – 600°C region. Above 600°C transformation of the film from an amorphous to a crystalline state takes place with a further decrease in index and high optical waveguide losses due to scattering from microcrystallites [19].

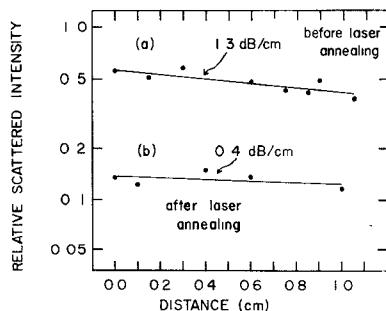


Fig. 4. Ta_2O_5 . Logarithm of the relative scattered intensity is plotted versus distance along the waveguide. The straight lines (a) and (b) represent best fits to the data obtained from sample Ta_2O_5 -2 before and after laser annealing, respectively.

The fact that laser annealing proved more effective on the directly sputtered Ta_2O_5 film than the thermally oxidized

film might be accounted for by the differences in film formation. The reactively sputtered oxide film, which was formed on an unheated substrate, should have had a more uniform amorphous composition than the thermally oxidized film. Laser annealing at a lower power density may have reduced bulk defects throughout the sputtered film without promoting the growth of microcrystals. In the case of the thermally oxidized film the higher power density used may have promoted microcrystalline growth while simultaneously annealing bulk defects.

To summarize the discussion of our results for the three waveguide materials, the Si₃N₄ waveguides were improved most effectively by CO₂ laser annealing. This is a reasonable result considering the strong absorption of Si₃N₄ at $\lambda = 10.6 \mu\text{m}$. A large improvement was noted after laser annealing a fair quality Nb₂O₅ waveguide, as was anticipated from the broad absorption band peaking at $\lambda = 11.8 \mu\text{m}$ for Nb₂O₅. Distinctly different results were obtained for the directly sputtered Ta₂O₅ waveguide and the thermally oxidized Ta₂O₅ waveguide. The infrared spectrum for Ta₂O₅ shows an absorption band peaking at $\lambda = 16.7 \mu\text{m}$ [23], suggesting that the absorption at 10.6 μm may not be very high. This may explain the smaller improvements in the Ta₂O₅ waveguides induced by laser annealing as compared with the improvements in the Si₃N₄ and Nb₂O₅ waveguides.

V. CONCLUSIONS

We have extended the technique of CO₂ laser annealing, which we have used previously to improve Corning 7059 glass and ZnO thin-film waveguides, to reduce optical scattering losses in Si₃N₄, Nb₂O₅, and Ta₂O₅ waveguides, all fabricated on thermally oxidized silicon substrates. Values of waveguide attenuation comparable with or better than the lowest losses previously reported for each type of waveguide were obtained for all three waveguide materials. The optimum power density for laser annealing was found to differ widely for the two Ta₂O₅ films which had been formed by different methods, with a larger improvement being obtained for the reactively sputtered film. These results have helped to establish CO₂ laser annealing as a powerful technique for reduction of optical waveguide scattering loss, representing a major advance in silicon-based integrated optics technology.

ACKNOWLEDGMENT

We wish to thank L. Grenon of the Motorola Semiconductor Group for supplying the Si₃N₄ waveguides.

REFERENCES

- [1] J. T. Boyd and D. B. Anderson, "Effect of waveguide optical scattering on the integrated optical spectrum analyzer dynamic range," *IEEE J. Quantum Electron.*, vol. QE-14, pp. 437-443, 1978.
- [2] S. Dutta, H. E. Jackson, and J. T. Boyd, "Reduction of scattering from a glass thin-film optical waveguide by CO₂ laser annealing," *Appl. Phys. Lett.*, vol. 37, pp. 512-514, 1980.
- [3] —, "Extremely low-loss glass thin-film optical waveguides utilizing surface coating and laser annealing," *J. Appl. Phys.*, vol. 52, pp. 3873-3875, 1981.
- [4] S. Dutta, H. E. Jackson, J. T. Boyd, F. S. Hickernell, and R. L. Davis, "Scattering loss reduction in ZnO optical waveguides by laser annealing," *Appl. Phys. Lett.*, vol. 39, pp. 206-208, 1981.
- [5] M. Gottlieb, G. B. Brandt, and J. J. Conroy, "Out-of-plane scattering in optical waveguides," *IEEE Trans. Circuits Syst.*, vol. CAS-26, pp. 1029-1035, 1979.
- [6] D. J. Channin, J. M. Hammer, and M. T. Duffy, "Scattering in ZnO-sapphire optical waveguides," *Appl. Opt.*, vol. 14, pp. 923-930, 1975.
- [7] T. Shiosaki, S. Ohnishi, Y. Hirokawa, and A. Kawabata, "As-grown CVD ZnO optical waveguides on sapphire," *Appl. Phys. Lett.*, vol. 33, pp. 406-407, 1978.
- [8] T. Shiosaki, S. Ohnishi, and A. Kawabata, "Optical properties of single-crystalline ZnO film smoothly chemical-vapor deposited on intermediately sputtered thin ZnO film on sapphire," *J. Appl. Phys.*, vol. 50, pp. 3113-3117, 1979.
- [9] T. Mitsuyu, S. Ono, and K. Wasa, "Structures and SAW properties of RF-sputtered single-crystal films of ZnO on sapphire," *J. Appl. Phys.*, vol. 51, pp. 2464-2470, 1980.
- [10] F. S. Hickernell, "Low loss zinc oxide optical waveguides on amorphous substrates," in *Tech. Dig. Topical Meeting Integrated and Guided-Wave Opt.*, Incline Village, NV, 1980, pp. WB6/1-4.
- [11] W. Stutius and W. Streifer, "Silicon nitride films on silicon for optical waveguides," *Appl. Opt.*, vol. 16, pp. 3218-3222, 1977.
- [12] W. M. Paulson, F. S. Hickernell, and R. L. Davis, "Effects of deposition parameters on optical loss for RF-sputtered Ta₂O₅ and Si₃N₄ waveguides," *J. Vac. Sci. Technol.*, vol. 16, pp. 307-310, 1977.
- [13] S. J. Ingrey and W. D. Westwood, "Birefringent waveguides prepared by reactive sputtering of niobium in O₂-N₂ mixtures," *Appl. Opt.*, vol. 15, pp. 607-609, 1976.
- [14] R. L. Davis and F. S. Hickernell, "An IO spectrum analyzer with thin-film lenses," in *Proc. 3rd Int. Conf. Integrated Optics and Optical Fiber Commun.*, San Francisco, CA, 1981, pp. WE6/1-4.
- [15] D. H. Hensler, J. D. Cuthbert, R. J. Martin, and P. K. Tien, "Optical propagation in sheet and pattern generated films of Ta₂O₅," *Appl. Opt.*, vol. 10, pp. 1037-1042, 1971.
- [16] S. J. Ingrey, W. D. Westwood, Y. C. Cheng, and J. Wei, "Variable refractive index and birefringent waveguides by sputtering tantalum in O₂-N₂ mixtures," *Appl. Opt.*, vol. 14, pp. 2194-2198, 1975.
- [17] D. B. Anderson, J. T. Boyd, M. C. Hamilton, and R. R. August, "An integrated-optical approach to the Fourier transform," *IEEE J. Quantum Electron.*, vol. QE-13, pp. 268-275, 1977.
- [18] R. L. Davis, "RF sputtered thin film niobium pentoxide for integrated optics," presented at the Amer. Vac. Soc. Meeting, Scottsdale, AZ, 1981.
- [19] R. K. Hickernell, "High optical quality amorphous tantalum oxide formed by the oxidation of sputtered tantalum films," *Bull. Amer. Phys. Soc.*, vol. 26, pp. 393-394, 1981.
- [20] H. Margenau and G. M. Murphy, *The Mathematics of Physics and Chemistry*. New York: Van Nostrand, 1943.
- [21] R. Ulrich and R. Torge, "Measurement of thin film parameters with a prism coupler," *Appl. Opt.*, vol. 12, pp. 2901-2908, 1973.
- [22] Yu. N. Volgin, G. P. Dubrovskii, and Yu. T. Ukhonov, "Lattice vibration spectrum of silicon nitride," *Sov. Phys. Solid State*, vol. 17, pp. 1089-1091, 1975.
- [23] R. A. Nyquist and R. O. Kagel, *Infrared Spectra of Inorganic Compounds (3800-45 cm⁻¹)*. New York: Academic, 1971.



Subhadra Dutta (S'81) was born in Shillong, India, on February 12, 1956. She received the B.Sc. degree (Special) in physics in 1976 from the University of the West Indies, Jamaica, and the M.S. and Ph.D. degrees in physics in 1978 and 1981, respectively, from the University of Cincinnati, Cincinnati, OH.

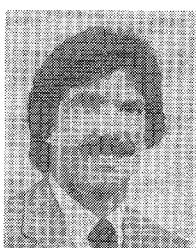
In her doctoral dissertation, she performed the first successful laser annealing experiments on thin film optical waveguides. After graduation, she joined the Westinghouse R&D Center, Pittsburgh, PA, where she is currently involved in research on the laser annealing of silicon.



Howard E. Jackson received the B.S. degree in physics from the University of Rochester, Rochester, NY, in 1965 and the Ph.D. degree in physics from Northwestern University, Evanston, IL, in 1971.

He was a Postdoctoral Research Fellow at the University of Toronto, Toronto, Ont., Canada, and McMaster University, Hamilton, Ont. His research interests include the low temperature thermal conductivity of insulators, heat pulses in solids, light scattering (including Raman scattering and Brillouin scattering) from both pure crystals and crystals with impurities, and from amorphous insulators. Presently, he is Associate Professor of Physics at the University of Cincinnati, Cincinnati, OH.

Dr. Jackson is a member of Sigma Xi and the American Physical Society.



Richard L. Davis was born in Southgate, CA, on January 13, 1947. He received the B.S. degree in physics from California State University, Long Beach, in 1972, and the M.S. degree in physics from California State University, Los Angeles, in 1973.

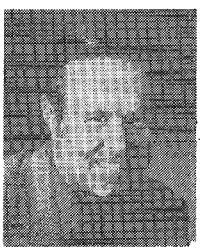
In 1973 he joined the Integrated and Non-linear Optics Group at Rockwell International, Anaheim, CA, where he conducted research on planar waveguide gas lasers and waveguide optics for coherent optical processing. In 1977 he joined the Government Electronics Division, Motorola, Inc., Scottsdale, AZ, where he directed the development of an integrated optic spectrum analyzer utilizing thin-film lenses. He is currently doing research on wide-band guided-wave Bragg cells.



J. T. Boyd (S'68-M'69-SM'78) was born in Martins Ferry, OH. He received the B.S. degree in electrical engineering from Pennsylvania State University, University Park, in 1964 and the M.S. and Ph.D. degrees in electrical engineering from Ohio State University, Columbus, in 1965 and 1969, respectively.

He is currently Professor of Electrical Engineering and Director of the Solid State Electronics Laboratories at the University of Cincinnati, Cincinnati, OH. He has conducted research in the areas of integrated optics and solid-state electronics. He has published over 70 papers in these areas. The research has been sponsored primarily by the Air Force Office of Scientific Research and the National Science Foundation. In support of this research he has received equipment grants from the National Science Foundation, numerous awards from the University of Cincinnati Research Council, and has been instrumental in securing a number of equipment donations from various industries. He has consulted for a number of companies since joining the University of Cincinnati in 1973. Prior to this he was employed by NCR and Rockwell International.

Dr. Boyd currently serves as Secretary of the Administrative Committee (AdCom) of the Quantum Electronics and Applications Society of the IEEE. He is a member of the Optical Society of America, Sigma Xi, Tau Beta Pi, and Eta Kappa Nu.



Fred S. Hickernell (M'64-SM'70) received the B.A. degree in education and the M.S. and Ph.D. degrees in physics from Arizona State University, Tempe, in 1953, 1959, and 1966, respectively.

Before 1960 he served as a Weather Officer in the U.S. Air Force and as an Engineer in the Theoretical Group of Goodyear Aircraft. In 1960 he joined Motorola Inc., Scottsdale, AZ, to work on the development of high-power microwave ferrite materials and devices. In 1963 he started work in the field of microwave acoustics, investigating the elastic and piezoelectric properties of III-V and II-VI compounds and their application to bulk acoustic mode traveling wave amplifiers and diffusion layer transducer delay lines. Since 1969 the acoustics work has been in the development of SAW components in standard, hybrid, and monolithic form for system applications. In 1978 the work was expanded to include integrated acousto-optical signal processing devices. He is presently a Principal Staff Engineer with Motorola's Government Electronics Division. He also serves as a Faculty Associate at Arizona State University, teaching undergraduate courses in the Mathematics Department.

Dr. Hickernell was elected to membership in Kappa Delta Pi, Chi Epsilon Pi, and Sigma Xi honor societies. He is active in the IEEE Group on Sonics and Ultrasonics and holds membership in the American Physical Society, the American Meteorological Society, and the American Scientific Affiliation. He has been honored by the IEEE, NASA, and Motorola for engineering contributions.